

# Lecture 1

# The Steady Magnetic Field

## Electromagnetic Field Theory



Sept. 25, 2019

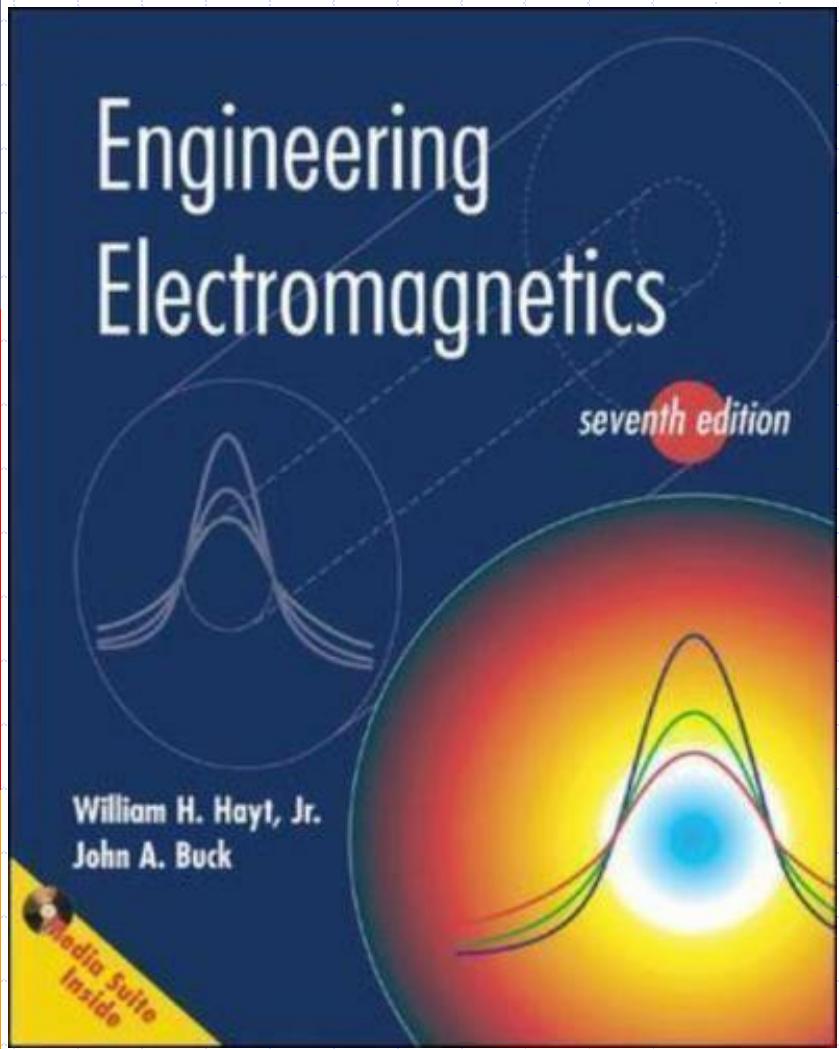
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# Selected Topics

- The steady magnetic field; Biot Savart and Ampere circuital laws
- Magnetic forces, torque, materials properties.
- Calculations of self and mutual Inductance
- Time varying field and Maxwell's equations
- The uniform plane wave; the transverse Electromagnetic (TEM) Wave
- Poynting theorem, normal and oblique incidence.



# Course Textbook



**W. Hayt Jr. and J. Buck,  
Engineering  
Electromagnetics,  
McGraw-Hill Education  
(India) Pvt Limited,  
2006.**



A very good book on  
electromagnetics. Easier to read  
and understand.  
Recommended



# Steady Magnetic Field

- The source of the steady magnetic field may be:
  - ➡ a permanent magnet  $\Rightarrow$  will be ignored
  - ➡ an electric field changing linearly with time  $\Rightarrow$  discussed later
  - ➡ a direct current.
- Our present study will concern the magnetic field produced by a *differential dc element* in free space.
- We may think of this *differential current element* as a vanishingly small section of a current-carrying filamentary conductor, where a filamentary conductor is the limiting case of a cylindrical conductor of circular cross section as the **radius approaches zero**.

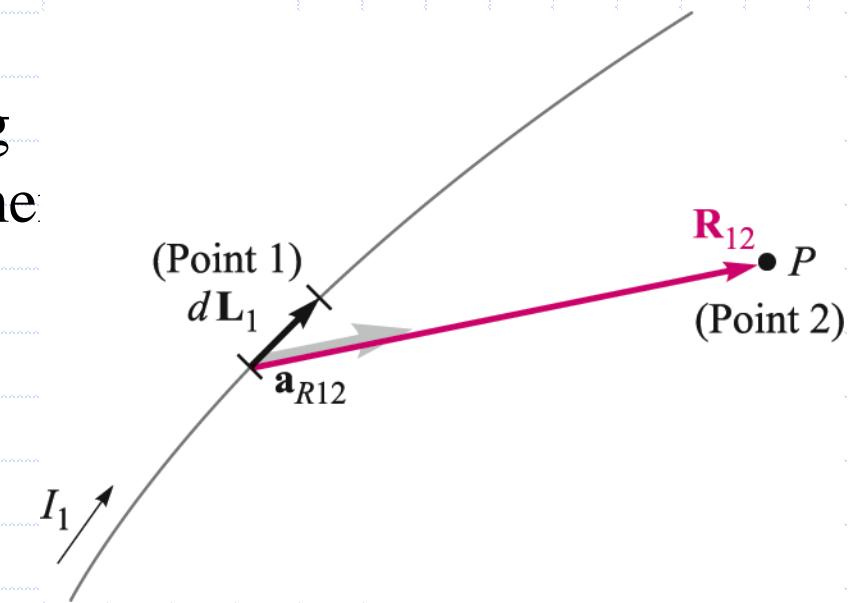
# Biot-Savart Law

The Biot-Savart Law specifies the magnetic field intensity,  $\mathbf{H}$ , arising from a “point source” current element of differential length  $d\mathbf{L}$ .

$$d\mathbf{H}_2 = \frac{I_1 d\mathbf{L}_1 \times \mathbf{a}_{R12}}{4\pi R_{12}^2}$$

The units of  $\mathbf{H}$  are [A/m]

Note in particular the inverse-square distance dependence, and the fact that the cross product will yield a field vector that points into the page. This is a formal statement of the right-hand rule



Note the similarity to Coulomb's Law, in which a point charge of magnitude  $dQ_1$  at Point 1 would generate electric field at Point 2 given by:

$$d\mathbf{E}_2 = \frac{dQ_1 \mathbf{a}_{R12}}{4\pi \epsilon_0 R_{12}^2}$$

# Magnetic Field Compared to Electric Field

## Distance

- ➡ The magnitude of the magnetic field varies as the **inverse** square of the distance from the source.
- ➡ The electric field due to a point charge also varies as the **inverse** square of the distance from the charge.

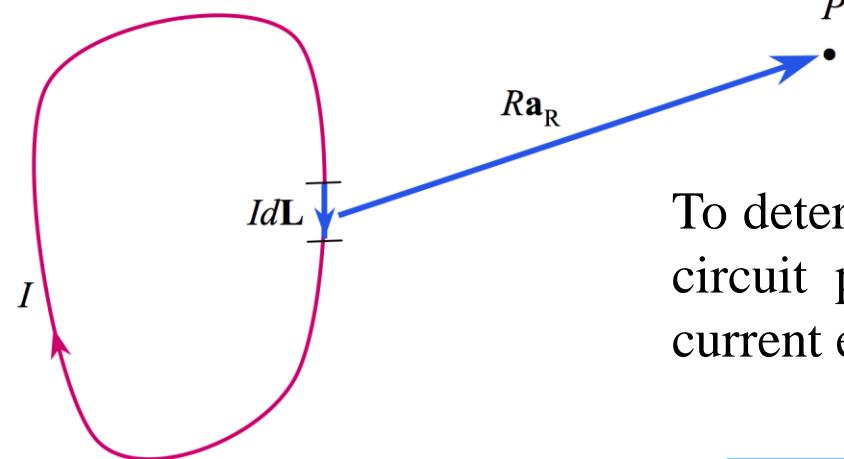
## Direction

- ➡ The electric field created by a point charge is **radial** in direction.
- ➡ The magnetic field created by a current element is **perpendicular** to both the length element  $d\mathbf{L}_1$  and the unit vector  $\mathbf{a}_R$ .

## Source

- ➡ An electric field is established by an **isolated electric charge**.
- ➡ The current element that produces a magnetic field must be **part of an extended current distribution**.
  - Therefore you must integrate over the entire current distribution.

# Magnetic Field Arising From a Circulating Current



At point  $P$ , the magnetic field associated with the differential current element  $Id\mathbf{L}$  is

$$d\mathbf{H} = \frac{Id\mathbf{L} \times \mathbf{a}_R}{4\pi R^2}$$

To determine the total field arising from the closed circuit path, we sum the contributions from the current elements that make up the entire loop, or

$$\mathbf{H} = \oint \frac{Id\mathbf{L} \times \mathbf{a}_R}{4\pi R^2}$$

Verified experimentally

The contribution to the field at  $P$  from any portion of the current will be just the above integral evaluated over just that portion.

# 2- and 3-Dimensional Currents

On a surface that carries uniform surface current density  $\mathbf{K}$  [A/m], the current within width  $b$  is

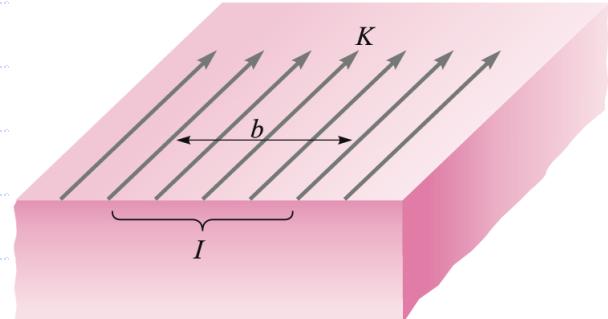
$$I = K b$$

.. and so the differential current quantity that appears in the Biot-Savart law becomes:

$$I d\mathbf{L} = \mathbf{K} dS$$

The magnetic field arising from a current sheet is thus found from the 2-D form of the Biot-Savart law:

In a similar way, a **volume current** will be made up of 3-D current elements, and so the Biot-Savart law for this case becomes:



$$\mathbf{H} = \int_s \frac{\mathbf{K} \times \mathbf{a}_R dS}{4\pi R^2}$$

$$\mathbf{H} = \int_{\text{vol}} \frac{\mathbf{J} \times \mathbf{a}_R dV}{4\pi R^2}$$

# Example of the Biot-Savart Law

In this example, we evaluate the magnetic field intensity on the  $y$  axis (equivalently in the  $xy$  plane) arising from a filament current of infinite length in on the  $z$  axis.

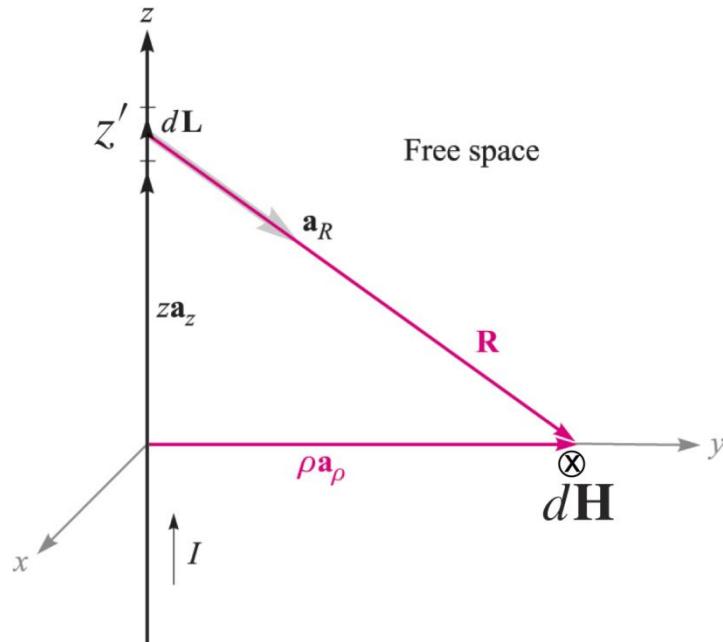
Using the drawing, we identify:

$$\mathbf{R} = \mathbf{r} - \mathbf{r}' = \rho \mathbf{a}_\rho - z' \mathbf{a}_z$$

and so..  $\mathbf{a}_R = \frac{\rho \mathbf{a}_\rho - z' \mathbf{a}_z}{\sqrt{\rho^2 + z'^2}}$

so that:

$$d\mathbf{H} = \frac{Id\mathbf{L} \times \mathbf{a}_R}{4\pi R^2} = \frac{Idz' \mathbf{a}_z \times (\rho \mathbf{a}_\rho - z' \mathbf{a}_z)}{4\pi(\rho^2 + z'^2)^{3/2}}$$



# Example: continued

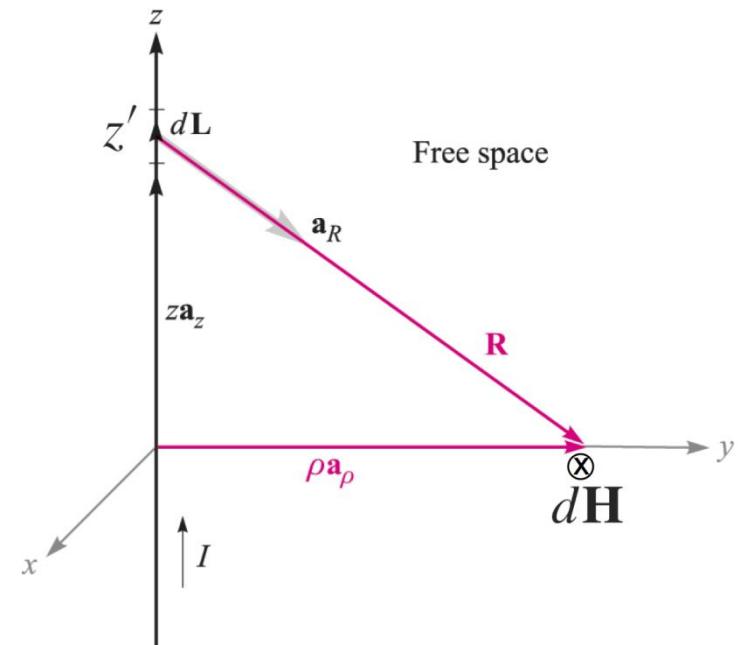
We now have:

$$d\mathbf{H} = \frac{Id\mathbf{L} \times \mathbf{a}_R}{4\pi R^2} = \frac{I dz' \mathbf{a}_z \times (\rho \mathbf{a}_\rho - z' \mathbf{a}_z)}{4\pi(\rho^2 + z'^2)^{3/2}}$$

Integrate this over the entire wire:

$$\mathbf{H}_2 = \int_{-\infty}^{\infty} \frac{I dz' \mathbf{a}_z \times (\rho \mathbf{a}_\rho - z' \mathbf{a}_z)}{4\pi(\rho^2 + z'^2)^{3/2}}$$

$$\mathbf{H} = \frac{I}{4\pi} \int_{-\infty}^{\infty} \frac{\rho dz' \mathbf{a}_\phi}{(\rho^2 + z'^2)^{3/2}}$$



..after carrying out the cross product

# Example: concluded

Evaluating the integral:

$$\int \frac{dx}{(a^2 + x^2)^{3/2}} = \frac{x}{a^2 \sqrt{a^2 + x^2}}$$

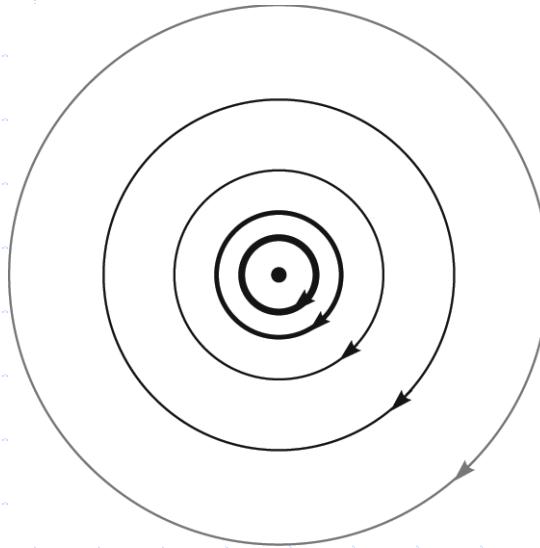
we have:  $\mathbf{H} = \frac{I}{4\pi} \int_{-\infty}^{\infty} \frac{\rho dz' \mathbf{a}_\phi}{(\rho^2 + z'^2)^{3/2}}$

$$= \frac{I \rho \mathbf{a}_\phi}{4\pi} \left. \frac{z'}{\rho^2 \sqrt{\rho^2 + z'^2}} \right|_{-\infty}^{\infty}$$

finally:

$$\mathbf{H}_2 = \frac{I}{2\pi\rho} \mathbf{a}_\phi$$

Current is into the page.  
 Magnetic field streamlines  
 are concentric circles, whose magnitudes  
 decrease as the inverse distance from the  $z$  axis



# Field Arising from a Finite Current Segment

In this case, the field is to be found in the  $xy$  plane at Point 2.

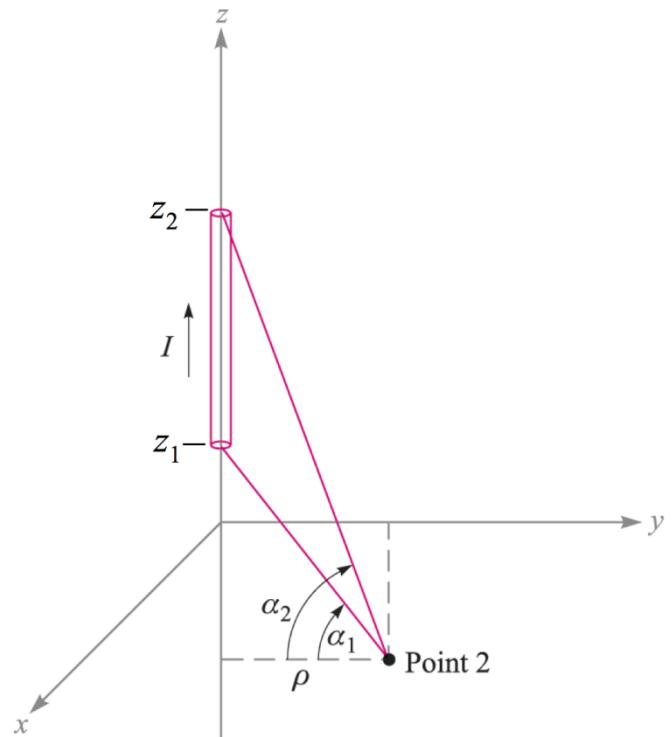
The Biot-Savart integral is taken **over the wire length**:

$$\mathbf{H} = \int_{z_1}^{z_2} \frac{Id\mathbf{L} \times \mathbf{a}_R}{4\pi R^2}$$

$$= \int_{\rho \tan \alpha_1}^{\rho \tan \alpha_2} \frac{Idz \mathbf{a}_z \times (\rho \mathbf{a}_\rho - z \mathbf{a}_z)}{4\pi(\rho^2 + z^2)^{3/2}}$$

..after a few additional steps, we find:

$$\mathbf{H} = \frac{I}{4\pi\rho} (\sin \alpha_2 - \sin \alpha_1) \mathbf{a}_\phi$$



$$\int \frac{dx}{(a^2 + x^2)^{3/2}} = \frac{x}{a^2 \sqrt{a^2 + x^2}}$$

# Another Example: Magnetic Field from a Current Loop

Consider a circular current loop of radius  $a$  in the  $x$ - $y$  plane, which carries steady current  $I$ . We wish to find the magnetic field strength anywhere on the  $z$  axis.

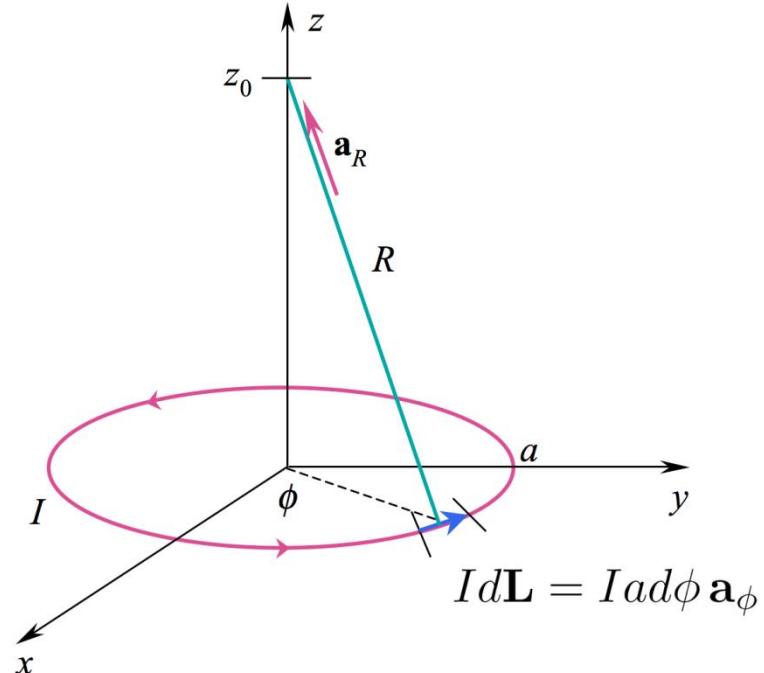
We will use the Biot-Savart Law:

$$\mathbf{H} = \int \frac{Id\mathbf{L} \times \mathbf{a}_R}{4\pi R^2}$$

where:  $Id\mathbf{L} = Iad\phi \mathbf{a}_\phi$

$$R = \sqrt{a^2 + z_0^2}$$

$$\mathbf{a}_R = \frac{z_0 \mathbf{a}_z - a \mathbf{a}_\rho}{\sqrt{a^2 + z_0^2}}$$



# Example: Continued

Substituting the previous expressions, the Biot-Savart Law becomes:

$$\mathbf{H} = \int_0^{2\pi} \frac{Iad\phi \mathbf{a}_\phi \times (z_0 \mathbf{a}_z - a \mathbf{a}_\rho)}{4\pi(a^2 + z_0^2)^{3/2}}$$

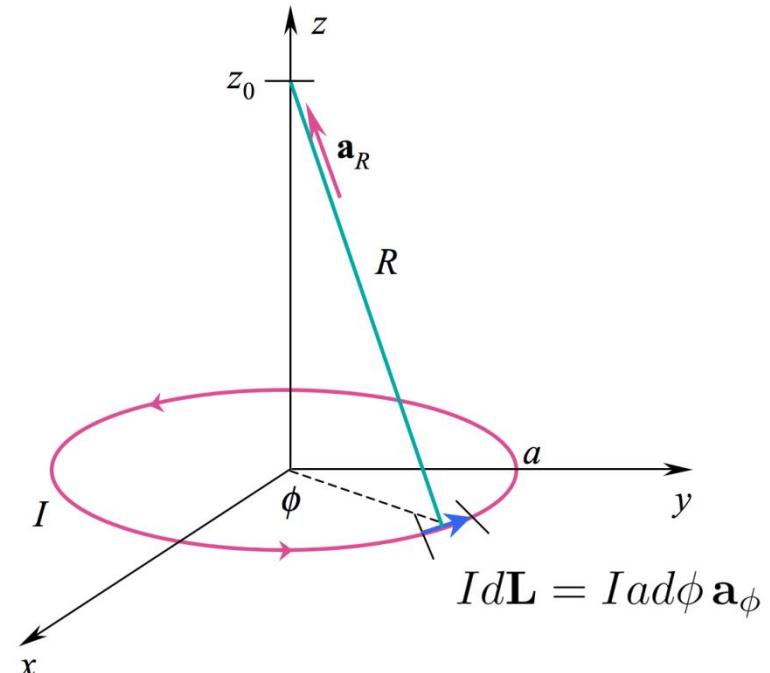
carry out the cross products to find:

$$\mathbf{H} = \int_0^{2\pi} \frac{Iad\phi (z_0 \mathbf{a}_\rho + a \mathbf{a}_z)}{4\pi(a^2 + z_0^2)^{3/2}}$$

but we must include the angle dependence in the radial unit vector:

$$\mathbf{a}_\rho = \cos \phi \mathbf{a}_x + \sin \phi \mathbf{a}_y$$

with this substitution, the radial component will integrate to zero, meaning that all radial components will **cancel on the z axis**.



# Example: Continued

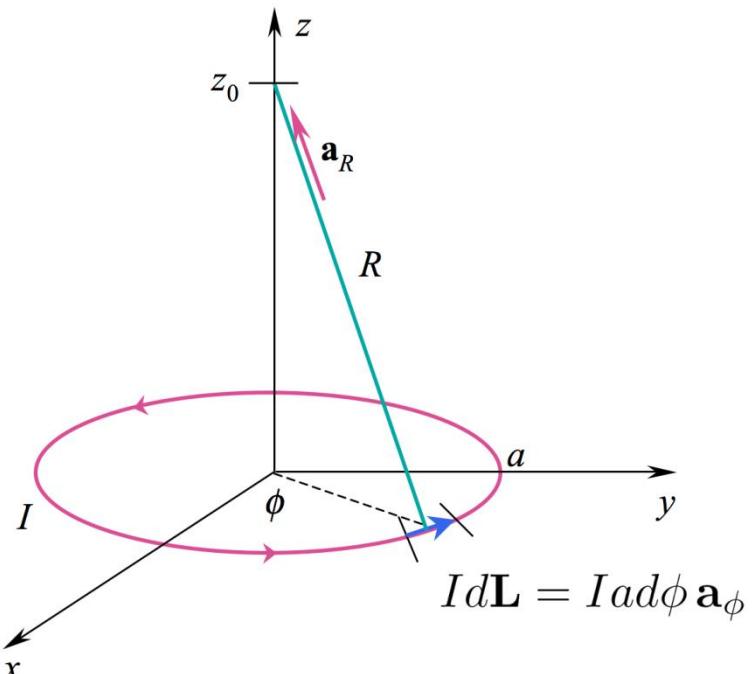
Now, only the  $z$  component remains, and the integral evaluates easily:

$$\mathbf{H} = \frac{I(\pi a^2) \mathbf{a}_z}{2\pi(a^2 + z_0^2)^{3/2}}$$

Note the form of the numerator: the product of the current and the loop area.

We define this as the *magnetic moment*:

$$\mathbf{m} = I(\pi a^2) \mathbf{a}_z$$



# Thank you



Have a nice day.

